

Measurement of bitumen viscosity in a room-temperature drop experiment: student education, public outreach and modern science in one

This content has been downloaded from IOPscience. Please scroll down to see the full text.

2014 Phys. Educ. 49 406

(<http://iopscience.iop.org/0031-9120/49/4/406>)

View [the table of contents for this issue](#), or go to the [journal homepage](#) for more

Download details:

This content was downloaded by: kostya_t

IP Address: 138.37.50.220

This content was downloaded on 16/07/2014 at 15:01

Please note that [terms and conditions apply](#).

Measurement of bitumen viscosity in a room-temperature drop experiment: student education, public outreach and modern science in one

A T Widdicombe¹, P Ravindrarajah¹, A Sapelkin¹, A E Phillips¹,
D Dunstan¹, M T Dove¹, V V Brazhkin² and K Trachenko¹

¹ School of Physics and Astronomy, Queen Mary University of London, Mile End Road, London, E1 4NS, UK

² Institute for High Pressure Physics, RAS, 142190, Moscow, Russia

E-mail: k.trachenko@qmul.ac.uk

Abstract

The slow flow of a viscous liquid is a thought-provoking experiment that challenges students, academics and the public to think about some fundamental questions in modern science. In the Queensland demonstration—the world's longest-running experiment, which has earned the Ig Nobel prize—one drop of pitch takes about ten years to fall, leading to problems for actually observing the drops. Here, we describe our recent demonstration of slowly-flowing bitumen where appreciable flow is observed on the timescale of months. The experiment is free from dissipative heating effects and has the potential to improve the accuracy of measurement. Bitumen viscosity was calculated by undergraduate students during the summer project. Worldwide access to the experiment, as it runs, is provided by webcams uploading the images to a dedicated website, enhancing the student education experience and promotion of science. This demonstration serves as an attractive student education exercise and stimulates the discussion of fundamental concepts and hotly debated ideas in modern physics research: the difference between solids and liquids, the nature of the liquid–glass transition, the emergence of long timescales in a physical process and the conflict between human intuition and physical reality.

1. Introduction

Of the three basic states of matter, liquids remain the least understood, compared to gases and solids. Strongly interacting disordered systems with dynamics intermediate between gases and solids [1], liquids continue to intrigue researchers. Surprising though it may seem, in this age of scientific advancement, we do not

understand even the most basic and fundamental properties of liquids, such as their specific heat. Landau and Lifshitz state twice in their statistical physics textbook [2] that liquid thermodynamic properties cannot be calculated in general form, in contrast to solids and gases. The stated reason is that interactions in a liquid are strong and system-specific, so that the energy and heat

Measurement of bitumen viscosity in a room-temperature drop experiment

capacity apparently become system-specific too. (Strong interactions are successfully treated in solids in the phonon approach, but this has been thought to be inapplicable to liquids where atomic displacements are large.) In addition to thermodynamic properties, strong interactions present a formidable challenge to calculating and understanding liquid transport properties such as viscosity [3, 4].

Viscosity is an important property of matter which measures its ability to flow under external perturbation. Most commonly discussed in liquids, viscosity has been studied in other states of matter too: in gases well over a century ago [5], in solids where flow is often referred to as ‘creep’ [4] and, more recently, in exotic states of matter such as quark–gluon plasma [6].

In liquids, measurable η changes with temperature in a very wide range, spanning 16 orders of magnitude: from $\eta \approx 10^{-3}$ Pa·s, as in low-viscous liquids such as room-temperature water or mercury, to about 10^{13} Pa·s at which point the liquid stops flowing at experimental timescales and forms glass by definition (if crystallization is avoided) [7]. Different types of experiment can be devised to measure η . The original falling ball Stokes experiment remains popular and able to uncover some non-trivial effects in liquids related to phase transformations [8].

In addition to being at the forefront of modern condensed matter physics research, issues surrounding viscosity and different ways to measure it have long been and remain of interest to physics educators [9–14] due to the wealth of physical concepts and effects that this property illustrates and highlights.

The experimental demonstration of slow pitch flow under gravity has been used for educational and science outreach purposes. Two notable examples include the experiment in the University of Queensland in Australia that started in 1927. This demonstration has been hailed the ‘world’s longest-running experiment’ and was awarded the Ig Nobel prize [15]. A similar demonstration started in Trinity College Dublin in 1944 [16]. In these experiments, drops of pitch form approximately every ten years [15–17]. These experiments have attracted the interest of major international news agencies, promoting and popularizing science, an exercise that has become increasingly important.

The main idea of the slow pitch flow demonstration is to confront long-time physical effects with human experience and intuition acquired at shorter times: pitch appears to be resistant to shear stress and immovable at short timescales and can even be shattered with a hammer, yet it flows at longer timescales as a liquid. The ability of the system to show elastic response at short time and viscous response at long time is often referred to as viscoelasticity, the property common to many systems, including familiar toys such as silly putty.

Apart from demonstration and outreach purposes, measurements of slow viscous flow have an important application. The ability to measure the viscosity of bitumen reliably has been discussed in [18]. In order to avoid dissipative heating present in common viscosity measurements, slow deformation of bitumen droplets was carefully measured. The measured viscosity was found to be an order of magnitude higher than in the previous measurements that involved dissipative heating [18].

Our rationale for installing a running bitumen experiment in the School of Physics and Astronomy in Queen Mary University of London was based on several ideas. First, stimulated by the media interest in our work on viscous flow and glass transition [19], we set out to make our idea visual, an important way to engage the students and enhance their learning experience as well as carry out a public outreach exercise.

Second, we sought to have a demonstration that challenges most fundamental concepts of our students, namely differences between states of matter and between liquids and solids in particular. The idea that the transition between the two phases can be continuous rather than discontinuous and that the difference between solid glasses and liquids is only quantitative but not qualitative is intriguing and has important scientific implications. The idea about the continuous nature of transition between solids and liquids was proposed by Frenkel a long time ago [20]. The idea was quickly criticised by Landau [21, 22] on the basis that the liquid–crystal transition involves symmetry changes and therefore cannot be continuous according to his theory of phase transitions. However, Frenkel was referring to solids in general, which include both crystals and glasses, emphasizing glass in his

later discussion [4]. Interestingly, the debate between Frenkel and Landau [20–22] took place right in between the launches of the two pitch flow experiments in Queensland and Dublin, suggesting that these questions were of increasing interest to both the experimental and theoretical communities at the time.

The discussion about whether the transition between liquids and solid is continuous or discontinuous is currently at the centre of the famous glass transition problem. According to the large set of experimental data, liquids and glasses are structurally identical [7] and the liquid–glass transition does not involve symmetry changes. Yet heat capacity at the liquid–glass transition changes with a jump, still suggesting a phase transition, possibly to a mysterious unknown glass phase [7]. This and related issues currently remain at the forefront of the problem of liquid–glass transition [7]. Considered by some to be the deepest and most interesting unsolved theoretical problem [23], the glass transition attracts recurring media attention (see, e.g., [19, 24]).

One does not need to assume the existence of a mysterious distinct glass phase and a phase transition in order to explain the heat capacity jump at the glass transition temperature T_g : the heat capacity jump is the inevitable result of liquid becoming too viscous to flow during the time of an experiment [25]. When the flow stops at the experimental timescale, a system's thermal and elastic properties change as a result, triggering the jump of heat capacity. In this picture, an equilibrium liquid above T_g becomes a non-equilibrium non-flowing liquid (solid glass) below T_g [25].

If the glass is just an extremely viscous liquid, can this non-intuitive concept be illustrated in a visual experiment? An obvious demonstration of this would be a very viscous and slow-flowing liquid that appears as a solid during short timescales (e.g. when an experimenter pushes and deforms the liquid with a pin), but starts flowing as a liquid at longer timescales: weeks, months and so on.

Third and importantly, we intended to observe and follow the bitumen flow at reasonably short timescales. Indeed, in the two pitch drop experiments discussed above, drops of pitch form approximately every ten years and were actually observed only recently when cameras recording the flow were installed and became operational [15, 16]. Our intention was to demonstrate noticeable bitumen flow (e.g. several grams) during the



Figure 1. Experimental setup showing top and bottom views of glass tubes with flowing bitumen. The radius of orifices increases from left to right. The thin bottom sections of the glass tubes have inner radii r and length l . The inner radius of the upper wide section of the tube is R .

timescale of one academic year. This way, new students starting in September could observe the

flow at the end of the final term in May next year. This applies to short-term visitors as well as associated students studying in the university for one year only.

2. Experiment and results

We now describe the design and installation of our experiment. The Total UK company provided two samples of bitumen of different (unspecified) viscosities (samples 35/50 and 95/25 with reported softening temperatures of 50–58 °C and 90–100 °C, respectively). We commissioned five glass tubes that narrow down at one end to thinner tubes (see figure 1) with different orifice radii $r = 2.5, 3.0, 4.0, 5.0$ and 6.0 mm. We did not have a reliable indication of the flow time originally, hence our reason for using different orifice radii was to have a wide range of flow times (flow time is proportional to r^4 , see below), increasing the possibility of observing the flow during a reasonable time window. As an additional way of potentially accelerating the flow, we prepared different weights to fit the glass tubes. The weights were not required in the actual experiment since the observable flow started approximately one month after the installation. The less viscous bitumen (sample 35/50 with reported softening temperature of 50–58 °C) was melted and kept in a furnace for 12 h at 90 °C. We used the less viscous sample in our attempt to observe the flow during the shorter timescale. The equilibrated bitumen melts were poured into the glass tubes in similar amounts. Rubber plugs were installed in the bottom thin sections to prevent the bitumen from pouring out. The plugs were removed after the bitumen was cooled to room temperature for 24 h and the flow started. During cooling, some of the plugs moved slightly, resulting in somewhat different initial levels of the bitumen in the thin bottom sections. During the experiment, room temperature variations were in the common 20–25 °C range.

Two webcams were installed, taking pictures every day and downloading them to the file repository connected to the dedicated website [26]. A computer program was written to update and show new photographs every day, copy them into a browsable gallery on the website and to join the photographs in the computer animation downloadable from the website. The website [26] has served

as an additional and important resource for student education as well as science demonstration and outreach. In the digital age, online demonstrations reporting experiments in real time such as this have the potential to add significantly to the way the education and promotion of science are carried out.

Bitumen viscosity η was calculated as part of the student summer project involving our first-year students, the first two authors of this paper. For a tube with radius r and length L , η and the flow rate are related by the Hagen–Poiseuille equation, the solution of the Navier–Stokes equation for the flow in the circular tube [27]:

$$\frac{dV}{dt} = \frac{P\pi r^4}{8\eta L}, \quad (1)$$

where V is volume, P is the pressure drop across the tube and $\frac{P}{L}$ is the constant pressure gradient.

We consider the flow from the thin bottom sections of the glass tubes with orifice radius r and length l in figure 1. The flow in the lower thin tube takes place under the hydrostatic pressure $g\rho h$, where h is the height of bitumen in the upper wide section of the tube and ρ is density. In addition, the flow takes place in the gravitational field, giving the extra term $g\rho$ added to the pressure gradient term ($\text{grad}P$) in the Euler and Navier–Stokes equations [27]. Consequently, the term representing the pressure gradient in the Hagen–Poiseuille equation, $\frac{P}{L}$, modifies as $\frac{g\rho h}{l} + g\rho$:

$$\frac{dV}{dt} = \frac{\pi r^4 g\rho(l+h)}{8\eta l}. \quad (2)$$

In the flowing bitumen experiment, equation (2) applies when P or h can be considered approximately constant. This is not the case in our experiment, where a considerable part of bitumen has flown out of the tubes (see figure 1). We therefore write $\frac{dV}{dt}$ as $\pi R^2 \frac{dh}{dt}$, where R is the radius of the upper wide section of the tube. This gives the first-order differential equation, integrating which gives:

$$\eta = \frac{g\rho r^4 t}{8R^2 l \ln \frac{h_1+l}{h_2+l}}, \quad (3)$$

where h_1 and h_2 are the heights of bitumen in the upper wide part of the glass tube at the beginning

and the end of the measurement, respectively, and t is time of the experiment.

Viscosity η was calculated on the basis of data collected in the flow experiment over $t = 317$ days. During this time period, the mass of bitumen that flowed and landed in the bottom flasks ranged from 5 g to 53 g for different tubes (the flow took place mostly in the form of almost non-interrupted strings of bitumen shown in figure 1 rather than in the form of identifiable drops). This provided enough flow to determine h_2 and h_1 with reasonable accuracy. We measured R , l and r using digital calipers with sub-millimeter precision.

Using the above data and equation (3), we calculated η in all five glass tubes. The averaged viscosity was found to be 7.5×10^6 Pa·s, with the relative error of 33% (the most substantial contribution to the error comes from non-systematic variations of h_2 due to different initial levels of bitumen in the thin bottom sections, as discussed above). Our bitumen sample is therefore about 30 times less viscous than in the world's longest (Queensland) experiment ($\eta = 2.3 \times 10^8$ Pa·s [15, 17]).

We note that our bitumen and that of Queensland are about 10^{10} – 10^{11} times more viscous than room-temperature water. At the same time, they are still 10^5 – 10^6 times less viscous than the system at the glass transition where $\eta \approx 10^{13}$ Pa·s, corresponding to the flow time exceeding the timescale of a typical experiment of the order of 10^3 s and 'apparent' solidification [7]. We comment on this further below.

3. Discussion and summary

We have described an experimental setup that allows students, researchers and the public to observe the flow of viscous bitumen during a time period of under one year. This timescale keeps a reasonable balance between fast and slow times: our bitumen appears solid on the timescale of days, yet shows appreciable flow over several months. By providing worldwide access to the experiment and transmitting it in real time, current information technology has the potential to enhance student education experience and public engagement. In view of the current interest in measuring viscosity in slow experiments without dissipative heating [18], the proposed

experimental setup may be used in further bitumen research and applications.

We make three comments regarding the implications of our demonstration for student education and research.

First, the bitumen we used is less viscous than in the previous demonstrations: in the Queensland and Dublin experiments, it takes about ten years for one drop to fall [15–17]. Public and media attention to these intriguing experiments was drawn primarily because this appears to be unusually long. It is important to remember that 'long' here refers to human timescales only. From the physical point of view, human timescales do not occupy a special place compared to others. Physical processes in Nature can involve vastly different timescales and liquid flow serves as a simple demonstration of this. Let us consider the property closely related to viscosity, the average time it takes an atom or molecule in liquid to jump between two adjacent positions [4]. If we melt a familiar SiO_2 system and cool it to room temperature without crystallization, this time is calculated to be 10^{67} s, 50 orders of magnitude longer than the age of the universe [25]. This time corresponds to viscosity of about 10^{77} Pa·s, over 70 orders of magnitude larger than the viscosity of our bitumen. It is in the appreciation of these scales of time and viscosity that the additional usefulness of our demonstration lies.

Second, our experiment highlights and contrasts human experience and intuition on the one hand and physical laws and processes on the other. A conflict between human experience/intuition and physical laws exists in several areas of physics. Quantum mechanics and general relativity are two notable examples where human experience is in striking conflict with physical laws at vastly different energy and length scales. Slowly flowing bitumen effectively highlights the same conflict, but in terms of time: our experience of the apparently solid bitumen, acquired at short observation time, contradicts the liquid behavior at longer times.

Third and finally, this demonstration raises an interesting question related to long timescales in physics: if the flow time becomes longer than any practical experiment and hence the experiment only probes the solid properties of the system, is it useful to debate whether the system is a slow-flowing liquid or a solid? This question brings us back to the original debate between Frenkel and Landau

[4, 20–22] and the continuing discussion about the nature of the liquid–glass transition [7, 23–25]. Our answer to this question is positively ‘yes’, because the answer has important consequences for the theories and models that we develop to describe the physical world. Namely, the experimental absence of structural and symmetry changes between liquid and the glass [7] implies that previously developed theories and models based on phase transition concepts do not apply. This stimulates the development of new theories and paradigms [25] and therefore drives the field forward. Regardless of what glass transition theories and models become accepted in the future, the simple and accessible demonstration of slow bitumen flow continues to challenge our students and ourselves.

Acknowledgments

We are grateful to the School of Physics and Astronomy and EPSRC for support, Total UK for providing bitumen samples, G Simpson for experimental assistance and B Still, A Owen, D Bolmatov P Micakovic and T Arter for their support with IT and the website.

Received 12 April 2014, in final form 28 April 2014,
accepted for publication 7 May 2014
doi:10.1088/0031-9120/49/4/406

References

- [1] Brazhkin V V and Trachenko K 2012 *Phys. Today* **65** 68
[2] Landau L D and Lifshitz E M 1964 *Statistical Physics* (Moscow: Nauka)

- [3] Chapman S and Cowling T G 1991 *Mathematical Theory of Non-Uniform Gases* (Cambridge: Cambridge University Press)
[4] Frenkel J 1947 *Kinetic Theory of Liquids* Fowler R *et al* (Oxford: Oxford University Press)
[5] Maxwell J C 1867 *Phil. Trans. R. Soc. Lond.* **157** 49
[6] Policastro G, Son D T and Starinets A O 2001 *Phys. Rev. Lett.* **87** 081601
[7] Dyre J C 2006 *Rev. Mod. Phys.* **78** 953
[8] Brazhkin V V *et al* 2008 *Phys. Rev. Lett.* **100** 145701
[9] Davenport T C 1968 *Phys. Educ.* **3** 139
[10] Wright P G 1977 *Phys. Educ.* **12** 323
[11] Bowlit C 1975 *Phys. Educ.* **10** 102
[12] Agrawal D C and Menon V J 1999 *Phys. Educ.* **34** 149
[13] Sabaz L 2006 *Phys. Educ.* **41** 112
[14] Houari A 2011 *Phys. Educ.* **46** 688
[15] Nadis S 2005 *Nature* **437** 938
[16] Johnston R 2013 *Nature News* doi: 10.1038/nature.2013.13418
[17] Edgeworth R, Dalton B J and Parnell T 1984 *Eur. J. Phys.* **5** 198–200
[18] Moran K and Yeung A 2004 *Can. J. Chem. Eng.* **82** 813
[19] Gilliland B 2011 Glass: the laziest liquid of them all? *Metro* www.cosmonline.co.uk/blog/2011/12/04/glass-laziest-liquid-them-all
[20] Frenkel J 1935 *Nature* **136** 167
[21] Landau L 1936 *Nature* **138** 840
[22] Landau L 1937 *Phys. Zs. Sowjet* **11** 26
Landau L 1937 *JETP* **7** 19
[23] Anderson P W 1995 *Science* **267** 1615
[24] Chang K 2008 The nature of glass remains anything but clear *New York Times*
[25] Trachenko K and Brazhkin V V 2011 *Phys. Rev. B* **83** 014201
[26] See <http://lazyliquids.com>
[27] Landau L D and Lifshitz E M 1988 *Fluid Mechanics* (Moscow: Nauka)